# Spectral Resolution of the Split EPR Signals Induced by Illumination at 5 K from the $S_1$ , $S_3$ , and $S_0$ States in Photosystem II<sup>†</sup>

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ABSTRACT: S-State-dependent split EPR signals that are induced by illumination at cryogenic temperatures (5 K) have been measured in spinach photosystem II without interference from the  $Y_D^{\bullet}$  radical in the g  $\sim$  2 region. This allows us to present the first decay-associated spectra for the split signals, which originate from the CaMn<sub>4</sub> cluster in magnetic interaction with a nearby radical, presumably Y<sub>Z</sub>\*. The three split EPR signals that were investigated, "Split  $S_1$ ", "Split  $S_3$ ", and Split  $S_0$ ", all exhibit spectral features at g $\sim 2.0$  together with surrounding characteristic peaks and troughs. From microwave relaxation studies we can reach conclusions about which parts of the complex spectra belong together. Our analysis strongly indicates that the wings and the middle part of the split spectrum are parts of the same signal, since their decay kinetics in the dark at 5 K and microwave relaxation behavior are indistinguishable. In addition, our decay-associated spectra indicate that the  $g \sim 2.0$  part of the "Split S<sub>1</sub>" EPR spectrum contains a contribution from magnetically uncoupled  $Y_{Z}^{\bullet}$  as judged from the g value and 22 G line width of the EPR signal. The g value, 2.0033-2.0040, suggests that the oxidation of  $Y_Z$  at 5 K results in a partially protonated radical. Irrespective of the S state, a small amount of a carotenoid or chlorophyll radical was formed by the illumination. However, this had relaxation and decay characteristics that clearly distinguish this radical from the split signal spectra. In this paper, we present the "clean" spectra from the low-temperature illumination-induced split EPR signals from higher plants, which will provide the basis for further simulation studies.

Photosystem II (PSII)<sup>1</sup> in the thylakoid membrane of higher plants, algae, and cyanobacteria (I) catalyzes the light-driven reduction of plastoquinone using electrons ultimately derived from water. The core of PSII is composed of the D1–D2 reaction center dimer that is surrounded by CP43, CP47, Cyt  $b_{559}$ , 11 small subunits, and three extrinsic proteins on the luminal side of the membrane (2-6). The photon energy is captured by pigments in antenna proteins and directed to primary electron donor  $P_{680}$ . Excitation of  $P_{680}$  is followed by the primary charge separation giving the

P<sub>680</sub><sup>+</sup>pheophytin<sup>-</sup> charge pair. Pheo<sup>-</sup> is rapidly oxidized by Q<sub>A</sub> and subsequently by electron transfer to Q<sub>B</sub>, which after two consecutive reductions and protonation dissociates from PSII.  $P_{680}^+$  is reduced by a nearby tyrosine residue,  $Y_Z$ , which is deprotonated to a nearby base (probably D1-His190) (7-10), producing the neutral Yz\* radical. Yz\* is reduced from the CaMn<sub>4</sub> cluster and ultimately from water. The catalytic site for water oxidation is composed of four Mn ions, one  $Ca^{2+}$ , and  $Cl^{-}$  as a cofactor (4, 11). During water oxidation, it cycles through five intermediate states,  $S_0-S_4$  (12). Molecular oxygen is released in the  $S_3 \rightarrow [S_4] \rightarrow S_0$ transition. The S<sub>1</sub> state is stable and dominates in the darkadapted material. S<sub>2</sub> and S<sub>3</sub> are metastable intermediate states that can be studied after one or two flashes given to a sample incubated in the dark. So is the most reduced state and can be produced by subjecting a dark-adapted sample to three flashes.

Under certain conditions, a pathway involving Cyt  $b_{559}$ , a Car, and a chlorophyll denoted Chl<sub>Z</sub> provides electrons to  $P_{680}^+$  (13-16). This also holds true for  $Y_D$ , which is situated on the D2 protein in a manner homologous to that of  $Y_Z$  (17).  $Y_D$  is not functional in linear electron transport but participates in charge equilibria with the CaMn<sub>4</sub> cluster (18, 19).  $Y_D^{\bullet}$  has a very stable radical g = 2.0046 EPR spectrum, which facilitates detailed spectroscopic investigations with many techniques (20).  $Y_Z$  is more difficult to study because of the fast decay and the magnetic interference from the close-lying CaMn<sub>4</sub> cluster in intact PSII. Therefore,  $Y_Z$  has

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 $<sup>^{\</sup>rm l}$  Abbreviations: Car, carotenoid; Chl, chlorophyll; Chlz, secondary chlorophyll electron donor to  $P_{680}^{+}$ ; Cyt  $b_{559}$ , cytochrome  $b_{559}$ ; CP43 and CP47, chlorophyll proteins of 43 and 47 kDa, respectively; D1 and D2, core subunits in PSII; DAD, 3,6-diaminodurene; DMSO, dimethyl sulfoxide; dpa, dipicolylamine; EPR, electron paramagnetic resonance; MES, 4-morpholinoethanesulfonic acid; ML, multiline; NIR, near-infrared; OEC, oxygen-evolving complex;  $P_{1/2}$ , microwave power to reach half-saturation;  $P_{680}$ , primary electron donor of PSII; Pheo, pheophytine;  $P_{\it PBQ}$ , phenyl-p-benzoquinone; PSII, photosystem II; Qa and QB, primary and secondary quinone acceptors, respectively, in PSII;  $Y_{\rm Z}$ , tyrosine 161 on the D1 subunit;  $Y_{\rm D}$ , tyrosine 161 on the D2 subunit.

mostly been studied in Mn-depleted PSII, where the radical lives for 150-700 ms compared to 0.03-1 ms in intact material (21, 22). However, these studies are of limited value in understanding molecular details of  $Y_Z$  in the intact system. More functional probes to  $Y_Z^{\bullet}$  are the so-called split EPR signals from the OEC. These were first found in modified PSII depleted of  $Ca^{2+}$  or  $Cl^-$  or displaced by various treatments (early findings reviewed in ref 23). By illumination of these samples at 0 °C, a split EPR signal was formed. This was 160 G wide (the width varied with treatment) and centered around g=2.0 (24-40) and arises from an  $S=\frac{1}{2}$  radical in magnetic interaction with the  $CaMn_4$  cluster. The most favored candidate for the radical is  $Y_Z^{\bullet}$  (30, 40-44), and the signal was attributed to the  $S_2Y_Z^{\bullet}$  redox state.

The S<sub>2</sub>Y<sub>Z</sub>• split signals in inhibited systems have provided much structural and functional information about Y<sub>Z</sub> and the CaMn<sub>4</sub> cluster. With respect to the functional OEC, more information is provided by recently discovered "split signals" that are induced by illumination at cryogenic temperatures (4-20 K) in the  $S_1$  (45-48) and  $S_0$  states (46, 48). Both the "Split S<sub>1</sub>" and "Split S<sub>0</sub>" signals can be induced by a single flash at 5 K and oscillate with the respective S state (46). The signals have preliminarily been assigned to magnetic interaction between Yz\* and the CaMn4 cluster in the respective S state (45-47, 49). A signal similar to the "Split S<sub>1</sub>" signal can also be formed by NIR illumination (700– 900 nm) at 5 K of centers in the S<sub>2</sub> state (50). Approximately 50% of the PSII in the  $S_0$  state and 40% in the  $S_1$  state have been found to give rise to the split signals in Thermosynechococcus elongatus on the basis of the amount of Q<sub>A</sub>formed simultaneously during the induction of the split signal (48). These numbers are consistent with numbers determined by other indirect methods (45, 46). Thus, the "split-forming pathway" is a dominating pathway during light-induced electron transfer in and around the OEC at cryogenic temperatures.

NIR illumination at 50 K of intact PSII in the  $S_3$  state results in formation of a split g=2 signal. It was proposed that the infrared illumination directly excites the CaMn<sub>4</sub> cluster in the  $S_3$  state to a potential sufficiently high to oxidize the nearby  $Y_Z$ , and the signal was attributed to magnetic interaction between a modified  $S_2$  state and, most probably,  $Y_Z^{\bullet}$  (51–53). We study a similar signal induced in a  $S_3$  state with visible light at 5 K.

In this paper, the split EPR signals, originating from the  $S_1$ ,  $S_3$ , and  $S_0$  states, are investigated. These are induced by weak illumination at 5 K of PSII prepared in the different S states with a different number of flashes. Most published spectra of the split signals are obscured by the large EPR signal from the oxidized  $Y_D^{\bullet}$  dominating at g=2.00. Therefore, this spectroscopically important region has often been discarded in published spectra. Two exceptions are the  $S_2Y_Z^{\bullet}$  split signal in  $Ca^{2+}$ -depleted PSII and the NIR-induced  $S_3$  signal that recently have been measured in a mutant lacking  $Y_D$  in T. elongatus (54).

Here we study PSII from higher plants (spinach), which is quite different from the cyanobacterial enzyme, in samples where  $Y_D^{\bullet}$  was chemically reduced prior to the EPR measurements. In this way, the spectra from the light-induced split signals could be determined without interference from  $Y_D^{\bullet}$ . The first decay-associated spectra for the "Split  $S_1$ ", "Split  $S_3$ ", and "Split  $S_0$ " signals are presented. This also

enables microwave relaxation studies of all three signals that allow us to draw conclusions about which parts of the complex spectra belong together.

## MATERIALS AND METHODS

*PSII Membrane Preparation.* PSII-enriched membranes [BBYs (55)] were prepared from hydroponically grown greenhouse spinach (*Spinacia oleracea*) with modifications described in ref 56. The storage buffer was 25 mM MES-NaOH (pH 6.1), 15 mM NaCl, 3 mM MgCl<sub>2</sub>, and 400 mM sucrose. Chl determinations were made according to the work of Arnon (57). EPR measurements were performed at 2–4 mg of Chl/mL.

Chemical Reduction of  $Y_D$ . To prevent EPR subtraction artifacts and a spectral contribution in the  $g \sim 2$  region from the large EPR spectrum from  $Y_D^{\bullet}$ ,  $Y_D^{\bullet}$  was reduced chemically by an ascorbate/DAD treatment in complete darkness (18, 58) to contain only 1.5%  $Y_D^{\bullet}$  (Figure 1A, inset). The PSII particles (at 1.0 mg of Chl/mL) were dark-adapted for 1 h on ice, incubated with a 10 mM ascorbate/3 mM DAD solution for 40 min at 20 °C, and thereafter washed three times in storage buffer at 2 °C to remove the chemical reductants.  $P_DPQ$  [1 mM, dissolved in DMSO, final DMSO concentration of 2% (v/v)] was added to the reduced PSII sample before the EPR tubes were filled in the dark and rapidly frozen. These dark-adapted samples were used to study PSII in the  $S_1$  state.

Steady-State Oxygen Evolution. Steady-state oxygen evolution in saturating light at 20 °C was assessed with a Clark-type electrode (Hansatech Instruments). The Chl concentration was 10  $\mu g$  of Chl/mL in a buffer with 25 mM MES-NaOH (pH 6.1), 15 mM NaCl, 3 mM MgCl<sub>2</sub>, and 400 mM sucrose. PpBQ (0.5 mM in DMSO) was used as an electron acceptor. The oxygen evolution was 400  $\pm$  50  $\mu mol$  of  $O_2$  (mg of Chl) $^{-1}$ h $^{-1}$  and remained constant after the reduction protocol and subsequent washings.

Flash Advancement of the S States. The dark-adapted EPR samples with  $Y_D$  reduced and containing 0.5 mM PpBQ were given two or three saturating flashes at 5 Hz from a Nd: YAG laser (6 ns, 532 nm, 400 mJ/pulse) at 0 °C. After the flashes, the samples were then frozen within 1-2 s. In the presence of reduced  $Y_D$ , the flash-induced turnover of the S states is rapidly desynchronized as observed previously (18) due to misses, the presence of reduced  $Y_D$ , and the lack of a synchronizing preflash (58). A two-flash sample prepared with this reduction and flash protocol was dominated by the  $S_3$  state, although it also contained a considerable fraction of centers in the  $S_2$  state (58). The sample given three flashes was estimated to contain  $\sim$ 50% PSII in the  $S_0$  state and  $\sim$ 40% in the  $S_3$  state, which is normal in this type of samples.

EPR Spectroscopy. Low-temperature EPR measurements were performed in a completely dark laboratory with a Bruker ELEXYS500E spectrometer using a SuperX EPR049 microwave bridge and a Bruker ST4102 standard cavity. The system was fitted with an Oxford Instruments cryostat and temperature controller. The split EPR signals were induced by illumination directly in the cavity at 5 K via a light guide. The light intensity, measured at the position of the sample, was 40 or 160 W/m² using white light filtered through a 5 cm thick CuSO<sub>4</sub> (aqueous) filter. This illumination regime was applied for 20 s to induce the split signals if not otherwise indicated in the text.

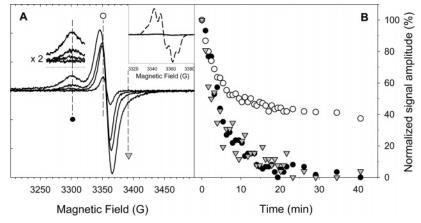


FIGURE 1: (A) Light minus dark difference spectrum of EPR signals induced by illumination at 5 K (10 s, 40 W/m<sup>2</sup>) of a dark-adapted PSII sample dominated by the S<sub>1</sub> state. The spectra were recorded immediately after the illumination at different microwave powers (from top to bottom at 3300 G: 25 mW, 1.0 mW, 0.1 mW, and 1.0 μW). The low-field peak region is enlarged 2 times for clarity. The dashed lines with symbols indicate the field positions used in Figure 1B. The inset shows the EPR spectrum of Y<sub>D</sub>\* before (dashed) and that remaining after the reduction treatment (black). (B) Decay in the dark at 5 K of the light-induced EPR spectrum recorded in the zero-flash sample at 25 mW, dominated by the "Split  $S_1$ " EPR signal. The decay is plotted at 3300 ( $\bullet$ ), 3350 ( $\bigcirc$ ), and 3390 G (gray triangles). The "Split S<sub>1</sub>" peak at 3300 G showed a fast, single decay phase ( $t_{1/2} = 3$  min) involving the entire amplitude of the signal. The same decay kinetics were found for the shoulder at 3390 G (gray triangles). At the peak of the signal (3350 G), 54% of the amplitude decayed with similar kinetics and the rest decayed much slower. EPR parameters for panels A and B; temperature, 5 K; microwave frequency, 9.41 GHz; modulation amplitude, 4 G. EPR parameters for the inset: temperature, 15 K; microwave power, 1.0 µW; modulation amplitude, 3.2 G.

To determine the spectral shapes of the split signals, decayassociated EPR spectra were constructed. The split signal in each respective S state was induced by the defined illumination protocol, and an EPR spectrum was recorded immediately after illumination. The decay in the dark of the EPR signals was then followed by recording a new spectrum every minute for  $\sim$ 45 min. The change in the signal intensity at every field position was plotted versus time and analyzed. In our analysis of the "Split S<sub>1</sub>" signal, a decay half-time of 3 min was found to encompass the total signal amplitude at the low-field peak at g = 2.035 (compare refs 46 and 48). The amplitude of this decay phase at each field position was set to represent the decay kinetics that belonged to the split signal part of the light-induced spectrum. A plot of this amplitude at all field positions produced the decay-associated spectrum of the "Split S<sub>1</sub>" signal. Similar protocols were applied to construct decay-associated spectra also for the signals originating from the S<sub>0</sub> state (in a three-flash sample) and the  $S_3$  state (in a two-flash sample).

The microwave power needed to reach half-saturation,  $P_{1/2}$ , for the different signals was determined graphically from a plot of  $\log(I/P^{0.5})$  versus  $\log P$  (milliwatts) (59, 60). I is the signal intensity at microwave power P.  $P_{1/2}$  was found by extrapolating the linear parts of the curve to their intersecting point, where the  $P_{1/2}$  value can be read on the abscissa directly.

Illumination of the sample for 3 min in room light (room temperature) followed by a 15 min dark adaptation provided complete oxidation of Y<sub>D</sub>\* (58). The EPR signal from completely oxidized Y<sub>D</sub>\* was used as an internal standard, and the signal was induced after all other measurements in each sample. At normal pH, YD is very stable, and after our pre-illumination and dark adaptation procedure, it is safe to assume that the reoxidized Y<sub>D</sub> amounts to one radical per PSII center (fully oxidized YD\* equals one spin per PSII).

## **RESULTS**

Spectral Resolution of the "Split  $S_1$ ", "Split  $S_3$ ", and "Split S<sub>0</sub>" EPR Signals. The light-induced EPR signals from PSII at 5 K are complex spectra with components from the split signals and light-induced Car/Chl radicals. Normally, the Y<sub>D</sub>\* radical EPR spectrum is very large and obscures the lightinduced signals in the  $g \sim 2.0$  region. To improve the spectral resolution of the light-induced EPR spectra, Y<sub>D</sub> was chemically reduced before the different S states were induced by laser flashes. In those samples, the split signals were induced with weak light at 5 K. All light-induced EPR signals were observed overlaid on a small remaining Y<sub>D</sub>• (1.5%, Figure 1A, inset), but in the figures, illuminated minus dark difference spectra are shown, which eliminate the contribution from the small YD.

Figure 1 describes the results for the "Split S<sub>1</sub>" signal that was studied in the sample incubated in the dark (zero flashes). Figure 1A shows the difference spectra recorded at different microwave powers, immediately after the 10 s illumination. At an elevated microwave power (>1 mW), the EPR spectrum is dominated by the typical shoulder at 3300 G (marked with a filled circle) known to originate from the "Split  $S_1$ " signal and a large radical-like feature at  $g \sim 2$ . At a very low microwave power (1  $\mu$ W), the shoulder at 3300 G is barely observable while the radical-like feature at  $g\sim 2$  remains. The microwave power dependence of the EPR spectrum is further described below. Figure 1B shows the decay in the dark of the EPR signal amplitude at different field positions recorded at a microwave power of 25 mW. At 3300 G (•), the entire signal decayed with a single exponential ( $t_{1/2} = 3$  min; Table 1). This reflects the "Split S<sub>1</sub>" decay, and the homogeneous decay kinetics indicates that the recorded spectrum is pure at this field position. The same monoexponential, complete decay was found on the shoulder at 3390 G (gray triangles), suggesting also a split signal component here. In contrast, at 3350 G (O), the fast

Table 1: Characteristics of the Light-Induced EPR Signals at 5 K

|                               | 0                   | 0                       |
|-------------------------------|---------------------|-------------------------|
| EPR signal                    | $t_{1/2}  (\min)^a$ | $P_{1/2}$ (mW)          |
| "Split S <sub>1</sub> "       | 3                   | 1                       |
| "Split S <sub>3</sub> "       | 3 (40%)             | 5                       |
| "Split S <sub>3</sub> "       | $slow^{b}$ (60%)    |                         |
| "Split S <sub>0</sub> "       | 3                   | 2                       |
| 9 G wide radical <sup>c</sup> | >30                 | $\sim 5 \times 10^{-3}$ |

 $^at_{1/2}$  for the decay in the dark at 5 K.  $^b$  Fraction of slowly or nondecaying "Split S<sub>3</sub>" signal.  $^c$  Radical, probably from a Car/Chl species, that is oxidized by our illumination regime in an S-state-independent way.

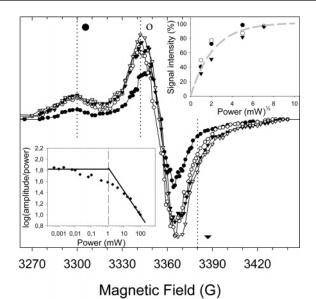


FIGURE 2: Decay-associated spectrum of the "Split S<sub>1</sub>" EPR signal at different microwave powers. The amplitude of the fast-decaying part ( $t_{1/2} = 3$  min) of the light-induced signal amplitude in the zeroflash sample, determined as described for Figure 1B, was used to construct the decay-associated spectrum of the "Split S<sub>1</sub>" EPR signal. The spectra were obtained at 1 ( $\bullet$ ), 4 ( $\bigcirc$ ), 25 ( $\blacktriangledown$ ), and 50 mW ( $\nabla$ ). The inset in the bottom left corner shows the microwave power dependence of the amplitude at 3300 G (the "Split S<sub>1</sub>" peak) plotted as log(amplitude/power<sup>1/2</sup>) vs power (milliwatts, log scale). The data are from the spectra in Figure 1A and from spectra recorded at additional microwave powers.  $P_{1/2}$  was determined to be 1 mW (59, 60). In the inset in the top right corner, power saturation data from three different parts of the created spectrum [3330 (●), 3342 (○), and 3380 G (▼)] are plotted and compared to a simulated curve with a  $P_{1/2}$  of 1 mW (gray dashed line). The dotted lines in the decay-associated spectrum indicate the field positions used to construct the inset. EPR parameters were as described for Figure 1.

decay ( $t_{1/2} = 3$  min) encompassed only 54% of the total EPR signal, while the rest of the signal decayed very slowly. Thus, at 3350 G, the recorded spectrum is mixed and composed of one part from the split signal and one part from another species. This other species could be a Car or Chl radical (see below).

From this kinetic analysis, a decay-associated spectrum was created between 3260 and 3450 G (Figure 2) for the fast ( $t_{1/2} \sim 3$  min) decaying signal. We propose that this spectrum reveals the "true" shape of the "Split S<sub>1</sub>" signal: a peak centered at 3300 G, a shoulder around 3390 G, and a 20 G wide radical-like feature centered at 3355 G.

The remaining, slowly decaying part in the  $g \sim 2.0$  region ( $t_{1/2} > 30$  min) was a simple radical signal (Figure 3), 9.5 G wide and centered at 3357.5 G ( $g = 2.0024 \pm 0.0001$ ). The

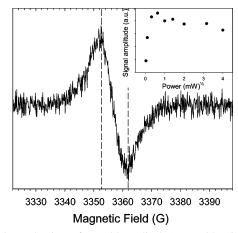


FIGURE 3: Induction of a stable radical  $(t_{1/2} > 30 \text{ min at 5 K})$  by a 20 s illumination (40 W/m²) at 5 K in a zero-flash sample dominated by the S<sub>1</sub> state. The spectrum is recorded after a 43 min decay in the dark at 5 K of the light-induced signal. The line width (dashed lines) is 9.5 G, and the signal is centered at g = 2.0024. The inset shows the microwave power dependence of this radical EPR signal. The power for half-saturation  $(P_{1/2})$  is  $\sim$ 5  $\mu$ W. EPR parameters: temperature, 5 K; microwave frequency, 9.41 GHz; microwave power, 1  $\mu$ W; modulation amplitude, 3.2 G.

radical was induced in  $\sim$ 5% of the PSII centers after illumination for 20 s at 40 W/m² as judged by double integration of the radical spectrum (compare also with the experiment described in Figure 4C; see below). When the light intensity is increased, more of this slowly decaying radical was formed, and a 20 s illumination at 160 W/m² gave rise to the radical in 10% of the PSII centers (not shown). The induction temperature (5 K), the width of this stable radical (9.5 G) (13), and the microwave power saturation at 5 K ( $P_{1/2} \sim 5 \mu$ W) (Figure 3, inset) indicated that this slowly decaying, light-induced species was a Car<sup>+</sup> radical, although an oxidized chlorophyll species cannot be ruled out.

The kinetic analysis of the illumination-induced species in the S<sub>1</sub> state (Figures 1 and 2) indicated that at least two spectral components were induced in the  $g \sim 2.0$  region of the spectrum. We therefore analyzed this part of the spectrum in further detail. Panels A and B of Figure 4 show the decayassociated spectra, resolved in the  $g\sim 2.0$  region, at microwave powers of 25 mW and 1  $\mu$ W. Figure 4C shows the fraction of PSII involved in the fast-decaying (ca. 4% of PSII) and slow-decaying or nondecaying parts (ca. 5% of PSII) of the spectra. At a low microwave power (1.0  $\mu$ W, Figure 4B,C), the spectrum representing the decaying part of the signal (•) was almost equal in size to the spectrum from the nondecaying part of the spectrum (O). At this low microwave power, both EPR spectra were recorded at or close to nonsaturating conditions, and the result indicates that approximately the same amount of spins was involved in each species. At a high microwave power (25 mW, Figure 4A), the amplitudes of the spectra were clearly different, and the decaying spectrum (•) was 2.7 times larger than the slow-decaying and nondecaying parts of the spectrum (O). Since both spectra represented approximately the same number of spins, the result indicates that the decaying part of the signal relaxes much faster than the nondecaying part of the spectrum. At a microwave power of 25 mW (Figure 4A), the nondecaying signal was severely saturated with

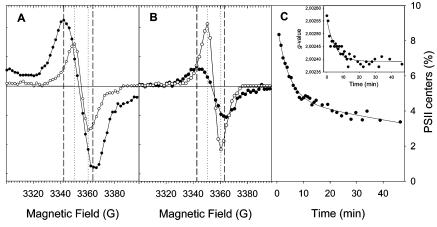


FIGURE 4: Decay-associated spectra in the  $g \sim 2.0$  region recorded at (A) 25 mW and (B) 1.0  $\mu$ W microwave powers after a 10 s illumination (40 W/m²) at 5 K of a dark-adapted sample dominated by the S<sub>1</sub> state. The spectra represent the amplitude of the light-induced spectrum decaying with a  $t_{1/2}$  of  $\sim 3$  min ( $\bullet$ ) and the nondecaying fraction of the light-induced spectrum ( $\circ$ ); compare Figure 3). The decaying spectrum had a line width of ca. 22 G (dashed lines), and the nondecaying spectrum had a line width of ca. 9.5 G (dotted lines). (C) Time-dependent decay of the radical feature measured on a PSII ( $Y_D^{\bullet}$  total) basis during incubation in the dark after illumination at 5 K. The inset shows the time-dependent decrease in the g value of the spectral feature at  $g \sim 2.0$  during the decay of the light-induced EPR signal. The spectra in panel C were recorded at a microwave power of 1.0  $\mu$ W. EPR parameters: temperature, 5 K; microwave frequency, 9.41 GHz; modulation amplitude, 4 G.

microwaves (compare Figure 3, inset) while the decaying signal was only saturated to a minor extent (compare Figure 2, bottom inset), resulting in the relative amplitudes observed in Figure 4A.

We can thus conclude that the two spectral components represent totally different species that both are induced by the illumination at 5 K. This conclusion is further substantiated by the g values and line widths of the decay-associated spectra in Figure 4. The nondecaying part of the spectrum ( $\bigcirc$ ) was characterized by g values of 2.0026 and 2.0024 in the high- and low-power spectra, respectively, similar to the g value of, for example, Car and Chl radicals. In contrast, the decaying part of the spectrum ( $\bigcirc$ ) was characterized by g values of 2.0040 and 2.0033 in the high- and low-power decay-associated spectra, respectively. These g values are significantly higher than the g value of the Car and Chl radicals and closer to the g value of, for example, a tyrosine radical (61, 62).

The difference between the g values found in the decay-associated spectra obtained at 25 mW and 1  $\mu$ W [ $t_{1/2} \sim 3$  min; Figure 4A,B ( $\bullet$ )] can be attributed to the possible contribution from a fast-decaying Car radical<sup>2</sup> (63) that will be mixed in the g value under the nonsaturating conditions (1  $\mu$ W, Figure 4B). In contrast at 25 mW, this Car<sub>fast</sub> will be severely saturated and will consequently contribute less to the g value of the decaying part of the spectrum [Figure 4A ( $\bullet$ )]. The slight difference in g value between the nondecaying spectra ( $\bigcirc$ ) at 25 mW and 1  $\mu$ W most probably reflects the fact that the spectrum obtained at 25 mW is distorted by the high microwave power due to rapid passage effects.

During the decay of the light-induced spectrum, this resulted in a decrease in the apparent g value due to the

progressively increasing domination of the nondecaying species (which exhibited a lower g value). This shift in the g value could be directly observed in the spectra recorded during the decay of the light-induced signal, and the time dependence of the g shift (recorded at a microwave power of 1  $\mu$ W) is shown in the inset of Figure 4C. In the first spectrum, recorded after the 10 s illumination, the g value of the central feature was 2.0026. At this stage, the g value was a mix (ca. 50% of each, Figure 4C) of the g values for the decaying (radical involved in the split signal and possibly Car<sub>fast</sub>) and nondecaying species (Car/Chl). The g value decreased with time in the dark to a stable g value of 2.0024 obtained after ca. 15 min. The half-time for this shift in the g value was ca. 3 min, indicating that it was directly correlated with the decay of the split part of the light-induced spectrum. The g shift could also be observed at higher microwave powers (not shown).

Interestingly, the decay-associated spectra in panels A and B of Figure 4 indicate that the line width of the spectral species decaying with a  $t_{1/2}$  of  $\sim$ 3 min was ca. 22 G (dashed lines). The broad line width was observable at both high and low microwave powers. The decaying species was thus much wider than the nondecaying signal that was only ca. 9–10 G wide (dotted lines in panels A and B of Figure 4; compare also with Figure 3).

To obtain the unperturbed spectrum of the "Split S<sub>3</sub>" signal, the decay kinetics of the light-induced signals at 5 K in a sample given two flashes were analyzed. The sample was dominated by the S<sub>3</sub> state but also contained a fraction of centers in the S<sub>2</sub> state due to misses in the turnover of the OEC. However, only the S<sub>3</sub> state will contribute to the light-induced split spectrum, since the S<sub>2</sub> state seemingly does not give rise to its own light-induced signals with our illumination protocol at 5 K. Figure 5A shows the light-induced minus dark difference spectra recorded at different microwave powers. The difference spectrum recorded at a microwave power of >1 mW was characterized by a broad peak at the low-field side of g=2 and a trough with two peaks at high field. The spectrum also contained a radical-

<sup>&</sup>lt;sup>2</sup> Optical spectroscopy experiments at 2 K in intact PSII revealed that light-induced Car<sup>+</sup> (induced by an illumination regime similar to that applied in our EPR experiments) decayed partly with a fast decay phase ( $t_{1/2} \sim 5$  min) and partly with a slower decay phase ( $t_{1/2} > 40$  min) (F. Ho, J. L. Hughes, K. G. V. Havelius, S. Styring, and E. Krausz, unpublished data).

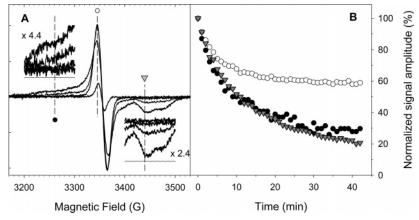


FIGURE 5: (A) Light minus dark difference spectra of EPR signals induced by illumination at 5 K (20 s, 40 W/m²) of PSII samples subjected to two laser flashes and dominated by the  $S_3$  state. The spectra were recorded immediately after the illumination at different microwave powers (from bottom to top at 3440 G: 32 mW, 1.0 mW, 0.1 mW, and 1.0  $\mu$ W). The peak region around 3260 G and the trough region around 3440 G are enlarged for clarity. The dashed lines with symbols at 3260 ( $\bullet$ ), 3345 ( $\bigcirc$ ), and 3440 G (gray triangles) indicate the field positions where the decay traces in panel B were measured. (B) Decay in the dark at 5 K of the light-induced (20 s, 40 W/m²) EPR spectrum recorded in the two-flash sample at a microwave power of 25 mW. In the decay traces at 3260 ( $\bullet$ ) and 3440 G (gray triangles), the fast-decaying phase with a  $t_{1/2}$  of  $\sim$ 3 min encompassed 40 and 30% of the signal, respectively. A large part of the signal remained after 43 min. The trace measured at the middle peak [3345 G ( $\bigcirc$ )] was dominated by the stable part (59%) but also contained a species decaying with a  $t_{1/2}$  of  $\sim$ 3 min (32%). EPR parameters: temperature, 5 K; microwave frequency, 9.41 GHz; modulation amplitude, 10 G.

like feature centered at  $g \sim 2.0$  quite similar to the "Split" S<sub>1</sub>" signal. In contrast, at a very low microwave power, the low- and high-field parts of the signal were hardly observable. Instead, the spectrum was dominated by a small radical feature around  $g \sim 2.0$ , which we also here interpret to belong mostly to an oxidized Car or Chl species. Figure 5B shows decay traces recorded at a high microwave power (25 mW) at 3260 (●), 3345 (○), and 3440 G (gray triangles), chosen to illustrate the decay kinetics of the light-induced signals over 43 min in the dark. At 3260 and 3440 G, the signal decayed with a  $t_{1/2}$  of  $\sim 3$  min (ca. 40 and 30%, respectively, of the total amplitude) and slower/nondecaying phases. It is clear from these decay traces that the signal did not decay to completion over our studied time interval (0-43 min). However, a complete decay, if possible to achieve, would involve EPR measurements over several hours, which are difficult to record with precision. Therefore, we were unable to determine the true decay half-time(s) for the slower phase. At 3345 G,  $\sim$ 30% of the signal decayed with a  $t_{1/2}$  of  $\sim$ 3 min, while more than 60% of the signal was stable in the dark at 5 K in this time range. In this analysis, we have constructed decay-associated spectra for the fast-decaying part of the signal and the signal that remained in the spectrum after a 40 min decay period.

Figure 6A shows the decay-associated spectrum for the fast-decaying ( $t_{1/2} \sim 3$  min) fraction of the spectrum in the two-flash sample. The fast-decaying "Split S<sub>3</sub>" signal spectrum ( $t_{1/2} \sim 3$  min) consisted of a broad peak centered at 3240 G, a slightly asymmetric derivative shape feature at 3354 G, and two troughs at 3440 and 3470 G. The nondecaying spectrum (Figure 6B, amplitude remaining after 43 min) was dominated by the shape of the Car or Chl radical species. Some "split-shaped" features were still observed at both wings in the spectrum, reflecting the fraction of the split spectrum that did not decay during 43 min.

A similar analysis was performed for the "Split S<sub>0</sub>" EPR signal recorded at different microwave powers (Figure 7A). This was studied in a sample given three flashes that was dominated by the  $S_0$  state. Due to misses, the sample also

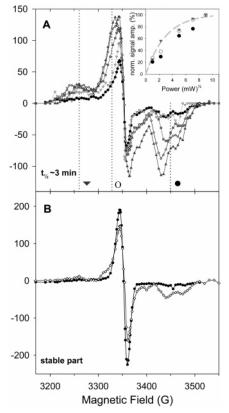


FIGURE 6: Decay-associated spectra of the "Split  $S_3$ " EPR signal. (A) The amplitude of the fast-decaying part  $(t_{1/2} \sim 3 \text{ min})$  of the light-induced signal amplitude in the two-flash sample, determined as described in the legend of Figure 5B, was used to construct the decay-associated spectrum of the "Split  $S_3$ " signal. The spectra were recorded at microwave powers of 1 ( $\blacksquare$ ), 5 (light gray triangles), 25 (gray squares), 50 ( $\diamondsuit$ ), and 81 mW (dark gray triangles). In the inset, power saturation data from three different parts of the created spectrum [3260 (gray triangles), 3335 ( $\bigcirc$ ), and 3450 G ( $\blacksquare$ )] are plotted and compared to a simulated curve with a  $P_{1/2}$  of 5 mW (gray dashed line). The dotted lines with symbols in the figure indicate the fields used in the inset. (B) Spectrum of the stable part of the signal remaining after a 43 min decay in the dark at 5 K in the two-flash sample created at microwave powers of 1 ( $\blacksquare$ ) and 50 mW ( $\diamondsuit$ ). EPR conditions were as described for Figure 5.

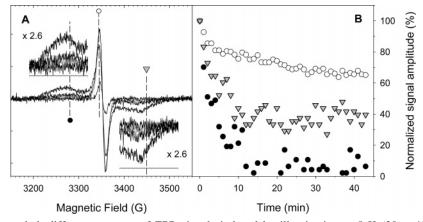


FIGURE 7: (A) Light minus dark difference spectra of EPR signals induced by illumination at 5 K (20 s, 40 W/m²) of PSII samples subjected to three laser flashes and dominated by the So state. The spectra were recorded immediately after the illumination at different microwave powers (from top to bottom at 3280 G: 25 mW, 1.0 mW, 0.2 mW, and 1.0  $\mu$ W). The dashed lines with symbols indicate the field positions where the decay traces in panel B were measured. (B) Decay in the dark at 5 K of the light-induced (20 s, 40 W/m<sup>2</sup>) EPR spectrum recorded in a three-flash sample. The decay traces were measured at 3280 (●), 3345 (○), and 3450 G (gray triangles). The peak at 3280 G ( $\bullet$ ) exhibited a fast single decay phase ( $t_{1/2} = 3$  min) involving the entire amplitude of the signal. At the trough at 3450 G (gray triangles), 62% of the signal amplitude decayed with these kinetics, while the rest remained stable over the measurement (43 min). At the middle peak, 3345 G (O), 18% of the signal amplitude decayed fast ( $t_{1/2} = 3$  min) and the rest much more slowly ( $t_{1/2} > 30$  min). EPR conditions were as described for Figure 5.

contained a significant fraction of centers in the S<sub>2</sub> and S<sub>3</sub> states. At microwave powers of > 1 mW, the low-field peak at 3280 G and a high-field trough at 3450 G dominated the spectrum together with the radical-like feature centered around  $g \sim 2.0$ . This spectrum is essentially similar to published EPR spectra of the "Split S<sub>0</sub>" signal (46, 48). At a low microwave power (1  $\mu$ W), the low- and high-field features were almost absent from the spectrum, which only showed the radical (most probably dominated by the Car and/or Chl radical) in the middle part of the spectrum. In essence, this behavior is similar to that for the "Split S<sub>1</sub>" and "Split S<sub>3</sub>" signals. Figure 7B shows the decay of the light-induced spectrum measured at different field positions. At the low-field peak, 3280 G (●), 100% of the light-induced signal decayed with a  $t_{1/2}$  of 3 min (Table 1), indicating that the spectrum was homogeneous at this field and that the entire "Split S<sub>0</sub>" signal decayed in this time range. At 3450 G (gray triangles), 62% of the signal was decaying with fast kinetics ( $t_{1/2} = 3$  min). The remaining 38% of the signal amplitude was not decaying in this time window (0-43 min), and we conclude that this fraction did not belong to the "Split  $S_0$ " signal. In the middle part of the signal [3340 G ( $\bigcirc$ )], only 18% of the light-induced signal decayed fast, while the rest was slowly decaying ( $t_{1/2} > 30 \text{ min}$ ).

The slowly/nondecaying part of the signal at 3450 G (triangles, Figure 7B) reflected that a fraction of PSII in the three-flash sample was in the S<sub>3</sub> state. The "Split S<sub>3</sub>" signal (Figures 5A and 6A) has a trough overlapping with the "Split S<sub>0</sub>" signal at this field position. As described above, the overlapping "Split S<sub>3</sub>" signal decayed both with fast ( $t_{1/2} \sim$ 3 min) and slower/nondecaying decay kinetics (Figure 5B). The slowly decaying species will not interfere with the decayassociated spectrum of the "Split S<sub>0</sub>" signal. However, the fast decay phase of the "Split S<sub>3</sub>" EPR signal will be included in the total fast decay kinetics that builds up the decayassociated spectrum of the three-flash sample. The middle feature that decayed fast (18%) belonged mainly to the S<sub>0</sub>state signal and to some extent to the fast-decaying part of the S<sub>3</sub>-state signal. However, only a minority of the PSII centers was in the S<sub>3</sub> state in the three-flash sample, and therefore, their contribution was quite small in the decaying part of the signal which was dominated by the "Split S<sub>0</sub>"

Figure 8 shows the spectrum associated with the  $t_{1/2} = 3$ min decay phase in the sample given three flashes. The signal is composed of a 140 G wide symmetric split signal (centered at 3355 G) and a 20 G wide radical-like feature around 3355 G. The latter is very similar to what was found for the signal in the S<sub>1</sub> state, while the 140 G wide symmetric part is similar to what was described earlier for the "Split S<sub>0</sub>" signal (Figure 7A) (46).

Microwave Power Saturation of the Split Signals. Fieldswept spectra recorded at different microwave powers (Figure 1A, "Split S<sub>1</sub>"; Figure 5A, "Split S<sub>3</sub>"; Figure 7A, "Split S<sub>0</sub>") allowed determination of the microwave power for halfsaturation,  $P_{1/2}$ , for the different signals. For the "Split S<sub>1</sub>" signal, the change in amplitude of the characteristic peak at 3300 G (g = 2.035) in a zero-flash sample was measured over a wide range of microwave powers, from 0.5  $\mu$ W to 100 mW (Figure 2, bottom inset). The microwave power for half-saturation,  $P_{1/2}$ , for the "Split S<sub>1</sub>" signal was 1 mW (Table 1). The microwave power saturation of the "Split  $S_0$ " EPR signal was monitored by the amplitude at the low-field peak, 3280 G, in a three-flash sample where the "Split S<sub>3</sub>" signal did not interfere. The  $P_{1/2}$  was 2 mW [Figure 8, bottom inset (♠); Table 1]. The microwave power saturation of the "Split S<sub>3</sub>" signal was measured at the trough at 3445 G in a two-flash sample [Figure 8, bottom inset  $(\diamondsuit)$ ]. The microwave power needed to saturate the "Split S<sub>3</sub>" signal was even higher, and  $P_{1/2}$  was 5 mW (Table 1). Thus, an unusually high microwave power was needed to saturate all three split signals compared to  $P_{1/2}$  values in the microwatt range for most organic radicals, including the Car/Chl species studied in these experiments (Figures 3 and 4). This indicates that very fast relaxing species give rise to the split signals. This implicates the involvement of a transition metal in all signals, indicating that the CaMn<sub>4</sub> cluster is part of the split signals (45, 46).

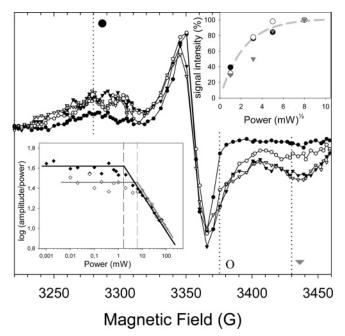


FIGURE 8: Decay-associated spectrum of the "Split S<sub>0</sub>" EPR signal. The amplitude of the fast-decaying part ( $t_{1/2} = 3 \text{ min}$ ) of the lightinduced signal amplitude in the three-flash sample, determined as described in the legend of Figure 7B, was used to construct the decay-associated spectrum of the "Split S<sub>0</sub>" signal at several microwave powers. The spectra were obtained at microwave powers of 1 ( $\bullet$ ), 10 ( $\circ$ ), 25 ( $\nabla$ ), and 64 mW ( $\circ$ ). The inset in the bottom left corner shows the microwave power dependence at 3280 G (◆) in a three-flash sample (the "Split  $S_0$ " peak) and at 3445 G  $(\diamondsuit)$  in a two-flash sample (the "Split  $S_3$ " trough).  $P_{1/2}$  was determined to be 2 and 5 mW for the "Split  $S_0$ " and "Split  $S_3$ " signals, respectively. In the inset in the top right corner, power saturation data from three different parts of the created spectrum [3280 (•), 3375 (O), and 3430 G (gray triangles)] are plotted and compared to a simulated curve with a  $P_{1/2}$  of 2 mW (gray dashed line). The dotted lines in the spectrum indicate the fields used in the inset. EPR conditions: temperature, 5 K; microwave frequency, 9.44 GHz; modulation amplitude, 10 G.

Microwave Power Saturation of the Decay-Associated Spectra. Decay-associated spectra of the "Split  $S_1$ " signal were created at several different microwave powers (Figure 2). The entire spectrum, illustrated by the field positions at 3300 ( $\bullet$ ), 3342 ( $\circ$ ), and 3380 G (gray triangles) in Figure 2 (top inset), showed a microwave power dependence similar to that of the  $P_{1/2}$  of the characteristic peak at 3300 G (Figure 2, top inset). This suggests that all parts of the signal originate from the same species.

The decay-associated spectrum of the "Split  $S_0$ " signal (see above) was constructed for several microwave powers. Figure 8 shows the spectra obtained at 1, 10, 25, and 64 mW. This allowed us to follow the relaxation of all parts of the signal in the decay-associated spectrum. This is illustrated at three field positions: 3280 ( $\bullet$ ), 3375 ( $\bigcirc$ ), and 3430 G (gray triangles) (Figure 8, top inset). The low-field side of the spectrum [Figure 8, top inset, 3280 G ( $\bullet$ )] relaxed coherently with a  $P_{1/2}$  of 2 mW [Figure 8, bottom inset ( $\bullet$ ) and top inset (gray dashed line)]. The high-field trough [Figure 8, top inset, 3430 G (gray triangles)] was somewhat more difficult to saturate with microwaves. However, the high-field trough is a mix of the  $S_0$ - and  $S_3$ -state split signals. Therefore, the faster relaxation of the "Split  $S_3$ " signal [ $P_{1/2} = 5$  mW; Figure 8, bottom inset ( $\diamondsuit$ )] interferes

significantly in the analysis and explains the difference in magnetic relaxation between the high-field and low-field sides in the decay-associated spectrum. Importantly, a part of the high-field side of the signal was not significantly mixed with the "Split S<sub>3</sub>" signal, and it relaxed with a  $P_{1/2}$  of 2 mW [Figure 8, top inset, 3375 G (O)]. To conclude, the relaxation of the >3400 G signal was closer to the  $P_{1/2}$  of the "Split S<sub>3</sub>" signal than the rest of the decay-associated spectrum of the three-flash sample, which relaxed with a  $P_{1/2}$  of 2 mW.

Decay-associated spectra of the "Split S<sub>3</sub>" signals (Figure 6) were created for several microwave powers: 1, 5, 25, 50, and 81 mW. The decay-associated spectrum based on the decay half-time  $(t_{1/2})$  of  $\sim 3$  min showed a similar relaxation pattern over the entire spectrum. This is illustrated in the inset of Figure 6A, where the signal amplitudes at three field positions [3260 (gray triangles), 3335 (O), and 3450 G (●)] were plotted versus the square root of the microwave power. The change in amplitude with microwave power was compared to the power saturation for the "Split S<sub>3</sub>" signal (gray dashed line) obtained in a separate experiment of the total light-induced amplitude at 3445 G (see above). We conclude from this comparison that all the features in the  $t_{1/2} \sim 3$  min decay-associated spectrum of the "Split S<sub>3</sub>" signal reflected the same very fast relaxing species. The stable part of the spectrum in the two-flash sample (Figure 6B) was dominated by the slowly relaxing Car/Chl radical ( $P_{1/2} \sim 5 \mu W$ ; Figure 3, inset; Table 1), while the small split wings displayed the hallmark of the fastrelaxing species.

#### DISCUSSION

In this paper, we have constructed decay-associated spectra reflecting the split EPR signals induced by illumination at 5 K of PSII positioned in the  $S_1$ ,  $S_3$ , and  $S_0$  states. Their microwave power saturation characteristics were determined and found to be consistent over the entire spectrum but different between the three split signals.

The OEC-related split EPR signals have been studied since 1989, when the first papers were published of a split EPR signal from S<sub>2</sub>Y<sub>Z</sub>• in Ca<sup>2+</sup>-depleted PSII (24, 25). In many of the previous publications, the middle part around  $g \sim 2.0$ has been excluded from the spectrum because the Y<sub>D</sub> radical dominates in this region. The consequence is that possible contributions to the split signal around  $g \sim 2.0$  have been cut out and ignored. We have overcome this problem by chemical reduction of the YD before measurements. This opened a spectral window in the  $g \sim 2$  region, since our illumination protocol did not oxidize Y<sub>D</sub> at 5 K (at pH 6.3).<sup>3</sup> Another obstacle to overcome before we could describe the true shape of the split signals was the fact that our illumination protocol induced an organic radical together with the split signal (independent of the S state). However, the decay kinetics of the split signals and the organic radical in the dark were different. Thereby, the fast-decaying split signals could be separated from the slowly decaying organic radical by decay-associated spectral analysis.

<sup>&</sup>lt;sup>3</sup> In Mn-depleted PSII, oxidation of  $Y_D$  at 15 K has been demonstrated, but this only occurs at high pH and is steered with a p $K_a$  of  $\sim$ 7.6 for a single protonatable group (64).

The "Split S<sub>1</sub>" EPR spectrum (Figure 2) ( $t_{1/2} = 3 \text{ min}$ ) was composed of one "split" part, which was ~80 G wide centered around 3342 G, and a more radical-like part, which was 22 G wide centered at 3355 G. These two parts could be separated neither by decay time nor by microwave power saturation, and we conclude that they together constitute the true experimental shape of the "Split S<sub>1</sub>" signal spectrum. The "Split S<sub>1</sub>" signal spectrum has been simulated at two microwave frequencies, X- and W-band (47). At X-band, the overall features in the simulated spectrum were similar to those of our experimental spectrum, including both the low-field peak and the central radical-like signal (Figure 2), which strengthens the validity of the simulation. In the theoretical analysis, the system involved magnetic interaction between the organic radical and both the ground and first excited spin states of the CaMn<sub>4</sub> cluster in the  $S_1$  state (S =0 and S = 1, respectively). The "Split  $S_1$ " signal originates from spin multiplet M2, which arises from the exchange coupling between the first excited S = 1 state of the Mn cluster and the  $S = \frac{1}{2}$  state of the interacting radical, presumably Yz\*. Thus, the "Split S1" signal has been linked to the first excited, S = 1, spin state of the CaMn<sub>4</sub> cluster (47). However, the theoretical model (47) does not exclude population also of the M1 ground state ( $S = \frac{1}{2}$ ) of the involved radical at a very low temperature. Therefore, the EPR spectrum might also contain a contribution from the radical species itself even if this might not be easily observed in the presence of the interaction signal.

In the absence of direct experimental proof, the radical involved in these light-induced split signals is thought to be  $Y_Z^{\bullet}$ . A strong argument for this assignment is the suitable, short distance (ca. 6 Å) between  $Y_Z$  and the CaMn<sub>4</sub> cluster, which is necessary to explain this kind of split EPR signal (2-6, 11), the spectral similarities with other split signals from the  $S_2Y_Z^{\bullet}$  state in inhibited PSII centers, and the ability to form a ML signal at 77 K from the split-containing samples under certain conditions (45).

The data here might provide further evidence for the assignment of Y<sub>Z</sub>• as the radical that is involved in the "Split S<sub>1</sub>" signal. Detailed investigation of the radical feature in the "Split S<sub>1</sub>" signal (Figure 4) showed that the signal was characterized by a relatively high g value (2.0033-2.0040) and a broad line width (22 G).  $P_{1/2}$  was determined to be  $\sim$ 1 mW at 5 K (Figure 2), which is unusually high for a free radical. Among the known radical species in PSII, Y<sub>D</sub>• in PSII centers in the S<sub>1</sub> state has a  $P_{1/2}$  of  $\sim$ 46  $\mu$ W at 8 K (60), while the Car/Chl radical studied here has a  $P_{1/2}$  of  $\sim$ 5  $\mu$ W (Table 1). In contrast,  $Y_Z^{\bullet}$  relaxes much faster due to cross relaxation with the close-lying CaMn<sub>4</sub> cluster. There are no relaxation measurements available for Yz at 5 K in the presence of a fully active CaMn<sub>4</sub> cluster. However, in Ca<sup>2+</sup>-depleted PSII centers,  $Y_Z^{\bullet}$  was found to have a  $P_{1/2}$  of > 50 mW while Y<sub>D</sub>• was a slower relaxer with a  $P_{1/2}$  of 8–10 mW at room temperature (41). Thus, the fast relaxation of the radical points to the involvement of Y<sub>Z</sub>. An analogous argument can be made for the line width. The signal was 22 G wide (the line width at half-height). This is very different from those of Car and Chl radicals, which are 9–11 G wide (65; compare Figure 3), while Y<sub>Z</sub>• has a line width of 19– 22 G (22, 65). Thus, the broad line width also suggests the direct involvement of a tyrosine radical, presumably Yz\* (due to the fast relaxation; see above), in the light-induced spectrum.

The assignment of the radical spectrum to Yz\* is further supported by the g value (2.0033-2.0040, Figure 4A,B) that was significantly higher than that for a Car radical (2.0024; 13) and closer to that of Y<sub>Z</sub>\* (2.0046; 22, 61, 65). However, the g value we obtained (2.0033-2.0040) is smaller than that normally reported for  $Y_Z^{\bullet}$  (2.0046). In this respect, it is interesting that the g value of a tyrosine radical is very dependent on the protonation state of the radical with g values ranging from 2.0032 for fully protonated tyrosine radicals (66) to the higher g values reported for the deprotonated tyrosine radicals in, for example, PSII (2.0046), ribonucleotide reductase (2.0050) (62, 67), and a synthetic Ru-Tyr molecule (2.0045) (68). The lower g value we observe could thus be interpreted to indicate that Yz is partially protonated when formed in intact PSII by illumination at a very low temperature (5 K).

Formation of a protonated tyrosine radical has precedence in the literature, and a g value of 2.0029 was found in a Ru-Tyr-dpa compound (69), where the phenolic proton was locked in a hydrogen bond. A very similar situation was also reported for the light-induced oxidation of  $Y_D^*$  in Mndepleted PSII at elevated pH (70). In this case, illumination at 1.8 K resulted in the formation of a partially protonated  $Y_D^*$  radical. This was recognized in high-field EPR by an unusually low  $g_x$  of 2.00643. Relaxation at 77 K resulted in a shift to 2.00756 typical for  $Y_D^*$  generated at physiological temperatures. This was interpreted to indicate that the  $Y_D^*$  radical induced at 1.8 K was involved in a hydrogen bond to the nearby D2-His190, but at the low temperature, the proton remained closer to  $Y_D$  than normal.

Our observation here of a relatively low g value also for the radical we assign to  $Y_Z^{\bullet}$  indicates a similar situation also around  $Y_Z$ . It has been suggested (46) that  $Y_Z$  is involved in a low-barrier hydrogen bond. Here we extend this conclusion and suggest that deprotonation along this bond at 5 K results in a radical in which the phenolic proton resides abnormally close to the phenolic oxygen. This suggestion implies a restricted proton movement in the hydrogen bond and is also indicative that  $Y_Z$  is protonated prior to oxidation, at least in the  $S_1$  state.

The observation of both a split signal and free Y<sub>Z</sub>• from the same redox state also has precedence in the literature. In Ca<sup>2+</sup>-depleted PSII, the state that gives rise to the split S<sub>2</sub>Y<sub>Z</sub>• magnetic interaction signal (always recorded at 5–15 K) has been suggested to contain a considerable fraction of free, magnetically uncoupled Yz\* when the EPR spectrum is recorded at room temperature (30, 36, 39, 41). Albeit disputed (71, 72), these results from many groups indicate that Yz can behave as a free form or as a form magnetically coupled to the same redox state of the CaMn<sub>4</sub> cluster. Such a situation would be consistent with Yz having the ability to populate either the M1 or the M2 state in the theoretical model for the "Split  $S_1$ " signal (47). Whether the population ratio follows a pure Bolzman distribution or if factors involving the molecular structure around Yz\* and the CaMn<sub>4</sub> cluster are involved cannot be decided at this stage.

The "Split S<sub>0</sub>" EPR spectrum (Figures 7 and 8) ( $t_{1/2} = 3$  min) was also composed of two "parts": the broad split part (140 G) and the narrow radical-like part ( $\sim$ 20 G) both centered at 3355 G that could not be separated by either

decay time or microwave power saturation. Thus, the 20 G radical-like feature of the "Split S<sub>0</sub>" signal behaves as a real part of the signal spectrum, videlicet, with the same decay kinetics and relaxation behavior as the broad split part. When the "Split S<sub>0</sub>" signal was discovered, it was immediately observed that the symmetry of the broad part of the signal resembled that of the earlier studied S2YZ\* split signals in inhibited PSII (46). There are several simulations of these signals (35, 37, 38, 40), and the signal shape has been found to be very dependent on the nature of the magnetic coupling (exchange and dipolar) and the distance between the interacting spin systems. However, neither of the simulations has predicted the existence of a narrow radical-like feature in the center of the split signal in the inhibited systems. It is therefore highly relevant that such a narrow radical part was observed in the S<sub>2</sub>Y<sub>Z</sub>• signal in acetate-treated PSII from a mutant lacking Y<sub>D</sub> (D2-Y160F) (54). As in our case, the mutant allowed studies of the split signals without the spectral overlap of Y<sub>D</sub>. In this work, the narrow central line was described for the first time in addition to the expected splitting of the S<sub>2</sub>Y<sub>Z</sub> signal in acetate-treated PSII. The authors pointed out that new simulation attempts were probably needed to explain the unexpected middle part of the spectrum. Our new data, which define the central component together with the broad splitting as the "Split  $S_0$ " signal, strengthen the conclusion that new simulations of the split signals must also explain radical-like features in the spectra similar to what has been achieved for the "Split S<sub>1</sub>" signal (47).

The decay-associated spectrum of the "Split S<sub>3</sub>" signal based on a  $t_{1/2}$  of  $\sim 3$  min from spinach PSII (Figure 6A) was very similar in shape to the IR-induced S<sub>3</sub> signal recorded in a mutant from T. elongatus lacking  $Y_D$  (54). The different parts of the fast-decaying signal ( $t_{1/2} \sim 3$  min) could not be separated by the microwave power saturation and behaved as one signal with a  $P_{1/2}$  of  $\sim 5$  mW (Figure 6A, inset). Both our spectrum and the spectrum from T. elongatus are characterized by a peak at 3260 G, two troughs at 3440 and 3470 G, and a narrow signal in the  $g \sim 2$  region of the spectrum. The narrow signal in the  $g\sim 2$  region had a somewhat broader peak and a smaller asymmetric trough around  $g \sim 2$  (Figure 6A and ref 54) compared to the radicallike features of the "Split  $S_1$ " and "Split  $S_0$ " signals. In the literature, the "Split S<sub>3</sub>" signal is induced by NIR illumination at 50 K (51, 54, 73). Here we have used light filtered through a 5 cm thick CuSO<sub>4</sub> (aqueous) filter, which eliminates the IR light. We induce the split signal at 5 K and measure it immediately, while in previous publications, the signal was induced at 50 K and then measured at 4 K. The time it takes to transfer the sample between 50 and 4 K is likely to involve the substantial decay of any decaying spectral parts and could give rise to a smaller signal size. This makes a straightforward comparison between the NIR-induced signal and our "Split S<sub>3</sub>" signal difficult.

All three split signals are difficult to saturate with microwaves; consequently, they reflect species that relax very fast. This has been found before (45, 46), but our analysis here demonstrates the same relaxation behavior for the entire

spectra. This is particularly important for the narrow middle part which earlier has been difficult to study due to overlaying signals from different radical species (YD\*, Car, and Chlz). However, these other radicals relax much slower than the split signals, which allowed studies of the latter at very high microwave powers. The fast relaxation behavior of the split signals indicates that very fast relaxers are involved in the signal and strongly implicates the involvement of the CaMn<sub>4</sub> cluster. This holds also for the S<sub>2</sub>Y<sub>2</sub>. split signal in inhibited PSII (as Ca<sup>2+</sup> or Cl<sup>-</sup> depleted) (24, 28, 40, 41). It is interesting to note that the relaxation behavior is S-state-dependent. At 5 K, the "Split S<sub>1</sub>" signal is easier to saturate than the "Split  $S_0$ " signal, which is easier to saturate than the "Split S<sub>3</sub>" signal (Table 1). S-Statedependent relaxation enhancement of the YD radical has previously been observed (60, 74). Due to cross relaxation between Y<sub>D</sub>• and the fast-relaxing CaMn<sub>4</sub> cluster, the microwave power saturation of Y<sub>D</sub>• in the presence of the different S states showed a clear difference between the S<sub>1</sub> state, with the lowest  $P_{1/2}$  (46  $\mu$ W at 8 K), and the other S states ( $P_{1/2} > 105 \,\mu\text{W}$  at 8 K). The relaxation enhancement was strongly temperature dependent, and the S2 and S3 states proved to have very similar relaxations. Since we have not detected any split signal from PSII in the S2 state, a full comparison with the earlier data cannot be made, but it is clear that the "Split S3" signals relaxed significantly faster than the "Split S<sub>1</sub>" signal. It is thus likely that the S-statedependent relaxation enhancement of the split signals and the Y<sub>D</sub>• radical has the same origin.

The slowly decaying radical we observed in low yield together with the split signals is presumably a carotenoid radical, but a chlorophyll radical cannot be excluded. The resolution in our spectrum might not be sufficiently good for the assignment of the radical to either Car or Chl. However, earlier studies on the light-induced radicals in intact PSII have revealed that a Car is the preferred donor at a very low induction temperature (20 K) (75). In Mn-depleted PSII, some studies observe only Car<sup>+</sup> formation upon illumination at 20 K (13), while many others find that illumination at  $\geq 6$  K oxidized both Car and Chl (14, 63). The different  $P_{680}^+$  electron donors available in our reduced samples at 5 K were Car/Chl, Cyt  $b_{559}$ , and the split donor. The  $Car^+/Chl^+$  and  $Cyt \ b_{559}^{ox}$  formed by the illumination could be quantified. Our short and weak illumination at 5 K (40 W/m<sup>2</sup>, 20 s) resulted in oxidation of Cyt  $b_{559}$  in ~8% of the PSII centers (not shown) and in the oxidation of ca. 5% Car/Chl. An increased level of illumination resulted in further oxidation of both species in addition to the formation of larger split signals. The maximum induction level of the "Split S<sub>1</sub>" signal (not shown) in this study was quite similar to that in refs 46 and 48. Zhang et al. (48) found 11% oxidation of Cyt b<sub>559</sub> and oxidation of 16% Car/Chl under their particular experimental conditions. By also following the acceptor side, they could assign ~40% of the PSII to form the split signal in the  $S_1$  state. In our PSII preparation from higher plants, we have not been able to investigate the acceptor side quantitatively due to the presence of the exogenous acceptor PpBQ, but it seems reasonable to assume the same order of magnitude as published earlier (48).

<sup>&</sup>lt;sup>4</sup> We have observed that the "Split S<sub>3</sub>" signal can be induced using monochromatic light at 532 nm (J.-H. Su, K. G. V. Havelius, F. Ho, F. Mamedov, and S. Styring, manuscript in preparation).

## **CONCLUSION**

The three split EPR signals that were investigated, "Split  $S_1$ ", "Split  $S_3$ ", and "Split  $S_0$ ", all exhibit spectral features in the  $g\sim 2$  region together with, for each signal, characteristic peaks and troughs. It would be very improbable that the wings and the middle part of the split spectra be derived from different signals, since their decay kinetics and relaxation behavior are identical. Therefore, these spectra can provide the basis for further simulation studies of the split signals in higher plants. In addition, it is highly likely that the EPR spectrum, at least in the  $S_1$  state, also contains a contribution from  $Y_Z^{\bullet}$ , spectroscopically distinguishable as a fast-relaxing radical spectrum resolved in the decay-associated spectrum. A small amount of Car/Chl was also oxidized in all cases, but this radical had clearly distinguishable relaxation and decay characteristics.

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